## **Energy and Power 4.4**

**Definition 4.27.** For a signal q(t), the instantaneous power p(t) dissipated in the 1- $\Omega$  resister is  $p_q(t) = |g(t)|^2$  regardless of whether g(t) represents a voltage or a current. To emphasize the fact that this power is based upon unity resistance, it is often referred to as the **normalized power**.

**Definition 4.28.** The total (normalized) *energy* of a signal g(t) is given by

$$E_g = \int_{-\infty}^{+\infty} p_g(t) dt = \int_{-\infty}^{+\infty} |g(t)|^2 dt = \lim_{T \to \infty} \int_{-T}^{T} |g(t)|^2 dt.$$

4.29. By the Parseval's theorem discussed in 2.39, we have

$$E_g = \int_{-\infty}^{\infty} |g(t)|^2 dt = \int_{-\infty}^{\infty} |G(f)|^2 df.$$

**Definition 4.30.** The average (normalized) **power** of a signal g(t) is given general formula by

$$P_g = \lim_{T \to \infty} \frac{1}{T} \int_{-T/2}^{T/2} |g(t)|^2 dt = \lim_{T \to \infty} \frac{1}{2T} \int_{-T}^{T} |g(t)|^2 dt = \left\langle \left| g(t) \right|^2 \right\rangle$$

**Definition 4.31.** To simplify the notation, there are two operators that used angle brackets to define two frequently-used integrals: Some properties for the (1) g(t) = 5 < < 2> = 5

(a) The "time-average" operator:

$$\langle g \rangle \equiv \langle g(t) \rangle \equiv \lim_{T \to \infty} \frac{1}{T} \int_{-T/2}^{T/2} g(t) dt = \lim_{T \to \infty} \frac{1}{2T} \int_{-T}^{T} g(t) dt \quad (46)$$

=a<91>+b<92>

When < 91,92> = 0,

we say g, and g2 are or the gonal.

(47)

(b) The **inner-product** operator:

$$\langle g_1, g_2 \rangle \equiv \langle g_1(t), g_2(t) \rangle = \int_{-\infty}^{\infty} g_1(t) g_2^*(t) dt$$

**4.32.** Using the above definition, we may write

•  $E_q = \langle g, g \rangle = \langle G, G \rangle$  where  $G = \mathcal{F} \{g\}$ •  $P_g = \left\langle |g|^2 \right\rangle$ 

• Parseval's theorem:  $\langle g_1, g_2 \rangle = \langle G_1, G_2 \rangle$ where  $G_1 = \mathcal{F} \{g_1\}$  and  $G_2 = \mathcal{F} \{g_2\}$ 

**4.33.** Time-Averaging over Periodic Signal: For **periodic** signal g(t) with period  $T_0$ , the time-average operation in (46) can be simplified to

$$(g) = \frac{1}{T_{0}} \iint_{g} g(t) dt$$
where the integration is performed over a period of  $g$ .  
Example 4.34.  $\langle \cos(2\pi f_{0}t + \theta) \rangle = \frac{1}{T_{0}} \int_{cos(2\pi f_{0}t + \theta)} dt = \frac{0}{0} = 0$ 

$$= \begin{cases} 0, & f_{0} \neq 0 \\ cos(2\pi f_{0}t + \theta) \rangle = \frac{1}{T_{0}} \int_{cos(2\pi f_{0}t + \theta)} dt = \frac{0}{0} = 0$$

$$= \int_{cos(2\pi f_{0}t + \theta)} \int_{cos(2\pi f_{0}t + \theta)} dt = \int_$$

Example 4.39. Suppose  $g(t) = 2e^{j6\pi t} + 3e^{j8\pi t}$ . Find  $P_g$ .  $P_g = 2^2 + 3^2 = 4 + 9 = 13$   $f_{i} = 3$   $f_{i} = 3$ Example 4.40. Suppose  $g(t) = 2e^{j6\pi t} + 3e^{j6\pi t}$ . Find  $P_g$ .  $g = 5e^{3e\pi t}$   $f_{i} = f_{2}$   $f_{i} = f_{2}$  $f_{i} = f_{$ 

**Example 4.41.** Suppose  $g(t) = \cos(2\pi f_0 t + \theta)$ . Find  $P_g$ .

Here, there are several ways to calculate  $P_g$ . We can simply use Example 4.35. Alternatively, we can first decompose the cosine into complex exponential functions using the Euler's formula:

$$(1) < |\cos(i\pi f_{o}t + \theta)|^{2} > = \frac{1}{2} \qquad P_{g} = |c_{1}|^{2} + |c_{s}|^{2} = \frac{1}{4} \times 1 + \frac{1}{4} \times 1 = \frac{2}{4} = \frac{1}{2}$$

$$(2) g(t) = \frac{1}{2} e^{j(2\pi f_{o}t + \theta)} + \frac{1}{2} e^{-j(2\pi f_{o}t + \theta)} = \frac{1}{2} e^{j\theta} e^{j(2\pi f_{o}t + \theta)} + \frac{1}{2} e^{-j\theta} e^{j(2\pi f_{o}t + \theta)} + \frac{1}{2} e^{-j\theta} e^{j(2\pi f_{o}t + \theta)} = \frac{1}{2} e^{-j\theta} e^{j(2\pi f_{o}t + \theta)} + \frac{1}{2} e^{-j\theta} e^{j(2\pi f_{o}t + \theta)} = \frac{1}{2} e^{-j\theta} e^{j(2\pi f_{o}t + \theta)} e^{j(2\pi f_{o}t + \theta)} = \frac{1}{2} e^{-j\theta} e^{j(2\pi f_{o}t + \theta)} e^{j(2\pi f_{o}t + \theta)} = \frac{1}{2} e^{-j\theta} e^{j(2\pi f_{o}t + \theta)} e^{j(2\pi f_{o}t + \theta)} e^{j(2\pi f_{o}t + \theta)} = \frac{1}{2} e^{-j\theta} e^{j(2\pi f_{o}t + \theta)} e^{j$$

**4.42.** The (average) power of a sinusoidal signal  $g(t) = A\cos(2\pi f_0 t + \theta)$  is

$$P_g = \begin{cases} \frac{1}{2} |A|^2, & f_0 \neq 0, \\ |A|^2 \cos^2\theta, & f_0 = 0. \end{cases}$$

This property means any sinusoid with nonzero frequency can be written in the form

$$g(t) = \sqrt{2P_g} \cos\left(2\pi f_0 t + \theta\right).$$

**4.43.** Extension of 4.42: Consider sinusoids  $A_k \cos(2\pi f_k t + \theta_k)$  whose frequencies are positive and distinct. The (average) power of their sum

$$g(t) = \sum_{k} A_k \cos\left(2\pi f_k t + \theta_k\right)$$
$$P_g = \frac{1}{2} \sum_{k} |A_k|^2.$$

is

Assume fo≠0. **Example 4.44.** Suppose  $g(t) = 2\cos\left(2\pi\sqrt{3}t\right) + 4\cos\left(2\pi\sqrt{5}t\right)$ . Find  $P_g$ .

$$P_g = \frac{2^2}{2} + \frac{4^2}{2} = \frac{1}{2}(4+16) = 10$$

**4.45.** For *periodic* signal g(t) with period  $T_0$ , there is also no need to carry out the limiting operation to find its (average) power  $P_g$ . We only need to find an average carried out over a single period:

$$P_{3} = \frac{1}{2} \int |g(t)|^{2} dt = \frac{1}{2} \int 1^{2} dt P_{g} = \frac{1}{T_{0}} \int_{T_{0}} |g(t)|^{2} dt.$$

(a) When the corresponding Fourier series expansion  $g(t) = \sum_{n=-\infty}^{\infty} c_n e^{jn\omega_0 t}$  is known,

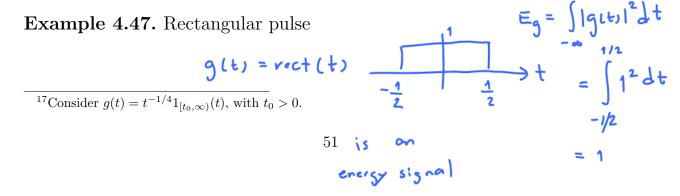
$$P_g = \sum_{k=-\infty}^{\infty} |c_k|^2$$

(b) When the signal g(t) is real-valued and its (compact) trigonometric Fourier series expansion  $g(t) = c_0 + 2 \sum_{k=1}^{\infty} |c_k| \cos(k\omega_0 t + \angle c_k)$  is known,

$$P_g = c_0^2 + 2\sum_{k=1}^{\infty} |c_k|^2$$

**Definition 4.46.** Based on Definitions 4.28 and 4.30, we can define three distinct classes of signals:

- (a) If  $E_g$  is finite and nonzero, g is referred to as an *energy signal*.
- (b) If  $P_g$  is finite and nonzero, g is referred to as a **power signal**.
- (c) Some signals<sup>17</sup> are neither energy nor power signals.
  - Note that the power signal has infinite energy and an energy signal has zero average power; thus the two categories are mutually exclusive.



Example 4.48. Sinc pulse 3-2 1 1 2 3 t

**Example 4.49.** For  $\alpha > 0$ ,  $g(t) = Ae^{-\alpha t} \mathbb{1}_{[0,\infty)}(t)$  is an energy signal with  $E_g = |A|^2/2\alpha$ .

 $E_g = \int |g(t)|^2 dt =$ 

 $|G(f)|^2 df = 1$ 

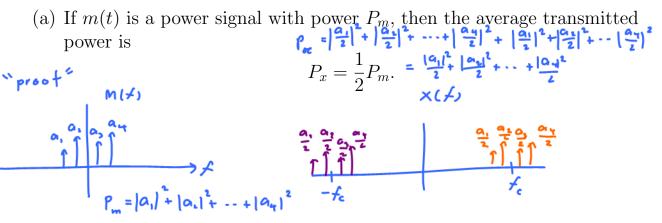
**Example 4.50.** The rotating phasor signal  $g(t) = Ae^{j(2\pi f_0 t + \theta)}$  is a power signal with  $P_g = |A|^2$ .

**Example 4.51.** The sinusoidal signal  $g(t) = A \cos(2\pi f_0 t + \theta)$  is a power signal with  $P_g = |A|^2/2$ .

4.52. Consider the transmitted signal

$$x(t) = m(t)\cos(2\pi f_c t + \theta)$$

in DSB-SC modulation. Suppose  $M(f - f_c)$  and  $M(f + f_c)$  do not overlap (in the frequency domain).



(b) If m(t) is an energy signal with energy  $E_m$ , then the transmitted energy is

$$E_x = \frac{1}{2}E_m$$

• Q: Why is the power (or energy) reduced?

## m(t) cus ( )

• Remark: When  $x(t) = \sqrt{2}m(t)\cos(2\pi f_c t + \theta)$  (with no overlapping between  $M(f - f_c)$  and  $M(f + f_c)$ ), we have  $P_x = P_m$ .